

# Erdős-Feller-Kolmogorov-Petrovsky law of the iterated logarithm for self-normalized martingales: a game-theoretic approach

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## Abstract

We prove an Erdős–Feller–Kolmogorov–Petrovsky law of the iterated logarithm for self-normalized martingales. Our proof is given in the framework of the game-theoretic probability of Shafer and Vovk. As many other game-theoretic proofs, our proof is self-contained and explicit.

*Keywords and phrases:* Bayesian strategy, constant-proportion betting strategy, lower class, upper class, self-normalized processes.

## 1 Main Result

Let  $S_n$  be a martingale with respect to a filtration  $\{\mathcal{F}_n\}_{n=0}^\infty$  and let  $x_n = S_n - S_{n-1}$  be the martingale difference. On some regularity conditions on the growth of  $|x_n|$ , various versions of the law of the iterated logarithm (LIL) have been given in literature. In particular the Erdős–Feller–Kolmogorov–Petrovsky law of the iterated logarithm (EFKP-LIL [16, Chapter 5.2]) is an important extension of LIL. Erdős [6] proved EFKP-LIL for symmetric Bernoulli random variables. EFKP-LIL has been generalized by Feller [7] for bounded and independent random variables and [8] (see also Bai [1]) for the i.i.d. case. Further, EFKP-LIL has been generalized for martingales by Strassen [19], Jain, Jogdeo and Stout [10], Philipp and Stout [15], Einmahl and Mason [5] and Berkes, Hörmann and Weber [2]. In particular, Einmahl and Mason [5] proved a martingale analogue of Feller’s result in [7], just as Stout [18] obtained a martingale analogue of Kolmogorov’s result in [11].

For self-normalized processes, EFKP-LIL was derived by [9, 3] in the i.i.d. case. However EFKP-LIL has not been derived in the martingale case, even though de la Peña, Klass and Lai [4] obtained the usual LIL. The purpose of this paper is to prove EFKP-LIL for self-normalized martingales. For a positive non-decreasing continuous function  $\psi(\lambda)$  let

$$I(\psi) := \int_1^\infty \psi(\lambda) e^{-\psi(\lambda)^2/2} \frac{d\lambda}{\lambda}. \quad (1)$$

We state our main theorem.

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**Theorem 1.1.** Let  $S_n$ ,  $n = 1, 2, \dots$ , be a martingale with  $S_0 = 0$  and  $x_n = S_n - S_{n-1}$  be a martingale difference with respect to a filtration  $\{\mathcal{F}_n\}_{n=0}^\infty$  such that

$$|x_n| \leq c_n \text{ a.s.}$$

for some  $\mathcal{F}_{n-1}$ -measurable random variable  $c_n$ . Let

$$A_n^2 := \sum_{i=1}^n x_i^2 \geq 0$$

and let  $\psi$  be a positive non-decreasing continuous function.

If  $I(\psi) < \infty$ , then

$$\mathbb{P}\left(S_n < A_n \psi(A_n^2) \text{ a.a.} \mid \lim A_n = \infty, \limsup c_n \frac{\psi(A_n^2)^3}{A_n} < \infty\right) = 1. \quad (2)$$

If  $I(\psi) = \infty$ , then

$$\mathbb{P}\left(S_n \geq A_n \psi(A_n^2) \text{ i.o.} \mid \lim A_n = \infty, \limsup c_n \frac{\psi(A_n^2)^3}{A_n} < \infty\right) = 1. \quad (3)$$

This theorem is a self-normalization of the result in Einmahl and Mason [5] and a generalization of the result in de la Peña, Klass and Lai [4]. The order of growth  $A_n/(\psi(A_n^2))^3$  for  $c_n$  is currently the best known order for EFKP-LIL even in the independent case ([2]). We call (2) the *validity* and (3) the *sharpness* of EFKP-LIL.

In (2) and (3), we are not assuming that the conditioning events happen with probability one. We can state (2) equivalently as

$$\mathbb{P}\left(\lim A_n = \infty, \limsup c_n \frac{\psi(A_n^2)^3}{A_n} < \infty, S_n \geq A_n \psi(A_n^2) \text{ i.o.}\right) = 0. \quad (4)$$

For our proof we adopt the framework of game-theoretic probability by Shafer and Vovk [17]. In a game-theoretic approach, for proving (2), we explicitly construct a non-negative martingale diverging to infinity on the event of (4).

We use the following notation throughout the paper

$$\ln_k n := \underbrace{\ln \ln \dots \ln n}_{k \text{ times}}$$

We also fix a small positive  $\delta$  for the rest of this paper, e.g.,  $\delta = 0.01$ . For our proof, as is often seen in the upper-lower class theory (cf. Feller [8, Lemma 1]), we can restrict our attention to  $\psi$  such that

$$\psi^L(n) \leq \psi(n) \leq \psi^U(n) \text{ for all sufficiently large } n, \quad (5)$$

where

$$\psi^L(n) := \sqrt{2 \ln_2 n + 3 \ln_3 n}, \quad \psi^U(n) := \sqrt{2 \ln_2 n + 4 \ln_3 n}.$$

Here  $L$  means the lower class and  $U$  means the upper class. It can be verified that  $I(\psi^U) < \infty$  and  $I(\psi^L) = \infty$ .

The rest of this paper is organized as follows. In Section 2 we give a game-theoretic statement corresponding to our main theorem. In Section 3 we give a proof of the validity and in Section 4 we give a proof of the sharpness.

## 2 Preliminaries on Game-Theoretic Probability

In order to state a game-theoretic version of Theorem 1.1, consider the following simplified predictably unbounded forecasting game (SPUFG, Section 5.1 of [17]) with the initial capital  $\alpha > 0$ .

SIMPLIFIED PREDICTABLY UNBOUNDED FORECASTING GAME

**Players:** Forecaster, Skeptic, Reality

**Protocol:**

$\mathcal{K}_0 := \alpha$ .

FOR  $n = 1, 2, \dots$ :

Forecaster announces  $c_n \geq 0$ .

Skeptic announces  $M_n \in \mathbb{R}$ .

Reality announces  $x_n \in [-c_n, c_n]$ .

$\mathcal{K}_n := \mathcal{K}_{n-1} + M_n x_n$ .

**Collateral Duties:** Skeptic must keep  $\mathcal{K}_n$  non-negative. Reality must keep  $\mathcal{K}_n$  from tending to infinity.

Usually  $\alpha$  is taken to be 1, but in Section 4 we use  $\alpha \neq 1$  for notational simplicity.

We prove the following theorem, which implies Theorem 1.1 by Chapter 8 of [17].

**Theorem 2.1.** *Consider SPUFG. Let  $\psi$  be a positive non-decreasing continuous function. If  $I(\psi) < \infty$ , Skeptic can force*

$$A_n^2 \rightarrow \infty \text{ and } \limsup c_n \frac{\psi(A_n^2)^3}{A_n} < \infty \Rightarrow S_n < A_n \psi(A_n^2) \text{ a.a.} \quad (6)$$

and if  $I(\psi) = \infty$ , Skeptic can force

$$A_n^2 \rightarrow \infty \text{ and } \limsup c_n \frac{\psi(A_n^2)^3}{A_n} < \infty \Rightarrow S_n \geq A_n \psi(A_n^2) \text{ i.o.} \quad (7)$$

We use the same line of arguments as in [14] and Chapter 5 of Shafer and Vovk [17]. We employ a Bayesian mixture of constant-proportion betting strategies. Here we give basic properties of constant-proportion betting strategies.

A constant-proportion betting strategy with betting proportion  $\gamma > 0$  sets

$$M_n = \gamma \mathcal{K}_{n-1}.$$

However,  $\mathcal{K}_n$  becomes negative if  $\gamma x_n < -1$ . For simplicity we consider applying the strategy (“keep the account open”) as long as  $\gamma c_n \leq \delta$  and sets  $M_n = 0$  once  $\gamma c_n > \delta$  happens (“freeze the account”). Define a stopping time

$$\sigma_\gamma := \min\{n \mid \gamma c_n > \delta\}. \quad (8)$$

Note the monotonicity of  $\sigma_\gamma$ , i.e.,  $\sigma_{\gamma'} \geq \sigma_\gamma$  if  $\gamma' \leq \gamma$ . We denote the capital process of the constant-proportion betting strategy with this stopping time by  $\mathcal{K}_n^\gamma$ . With the initial capital of  $\mathcal{K}_0^\gamma = \alpha$ , the value of  $\mathcal{K}_n^\gamma$  is written as

$$\mathcal{K}_n^\gamma = \alpha \prod_{i=1}^{\min(n, \sigma_\gamma - 1)} (1 + \gamma x_i).$$

By

$$t - \frac{t^2}{2} - t^2 \times |t| \leq \ln(1+t) \leq t - \frac{t^2}{2} + t^2 \times |t|$$

for  $|t| \leq \delta$ , taking the logarithm of  $\prod_{i=1}^n (1 + \gamma x_i)$ , for  $n < \sigma_\gamma$ , we have

$$\gamma S_n - \frac{\gamma^2 A_n^2}{2} - \gamma^3 A_n^2 \bar{c}_n \leq \ln(\mathcal{K}_n^\gamma / \alpha) \leq \gamma S_n - \frac{\gamma^2 A_n^2}{2} + \gamma^3 A_n^2 \bar{c}_n$$

and

$$e^{-\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2 / 2} \leq \mathcal{K}_n^\gamma / \alpha \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2 / 2}, \quad (9)$$

where

$$\bar{c}_n := \max_{1 \leq i \leq n} c_i.$$

We also set up some notation for expressing the condition in (6) and (7). An infinite sequence of Forecaster's and Reality's announces  $\omega = (c_1, x_1, c_2, x_2, \dots)$  is called a *path* and the set of paths  $\Omega = \{\omega\}$  is called the sample space. Define a subset  $\Omega_{<\infty}$  of  $\Omega$  as

$$\Omega_{<\infty} := \left\{ \omega \mid A_n^2 \rightarrow \infty, \limsup_n c_n \frac{\psi(A_n^2)^3}{A_n} < \infty \right\}.$$

For an arbitrary path  $\omega \in \Omega_{<\infty}$  we have

$$\exists C(\omega) < \infty, \exists n_1(\omega), \forall n > n_1(\omega), c_n < C(\omega) \frac{A_n}{\psi(A_n^2)^3}, \psi(A_n^2) \geq 1. \quad (10)$$

The last inequality holds by the lower bound in (5).

### 3 Validity

We prove the validity in (6) of Theorem 2.1. In this section we let  $\alpha = 1$ . We discretize the integral in (1) as

$$\sum_{k=1}^{\infty} \frac{\psi(k)}{k} e^{-\psi(k)^2/2} < \infty. \quad (11)$$

Since  $xe^{-x^2/2}$  is decreasing for  $x \geq 1$ , the function  $\lambda \mapsto \frac{\psi(\lambda)}{\lambda} e^{-\psi(\lambda)^2/2}$  is decreasing for  $\lambda$  such that  $\psi(\lambda) \geq 1$  and convergences of the integral in (1) and the sum in (11) are equivalent.

The convergence of the infinite series in (11) implies the existence of a non-decreasing sequence of positive reals  $a_k$  diverging to infinity ( $a_k \uparrow \infty$ ), such that the series multiplied term by term by  $a_k$  is still convergent:

$$Z := \sum_{k=1}^{\infty} a_k \frac{\psi(k)}{k} e^{-\psi(k)^2/2} < \infty.$$

This is easily seen by dividing the infinite series into blocks of sums less than or equal to  $1/2^k$  and multiplying the  $k$ -th block by  $k$  (see also [13, Lemma 4.15]).

For  $k \geq 1$  let

$$p_k := \frac{1}{Z} a_k \frac{\psi(k)}{k} e^{-\psi(k)^2/2}$$

and consider the capital process of a countable mixture of constant-proportion strategies

$$\mathcal{K}_n := \sum_{k=1}^{\infty} p_k \mathcal{K}_n^{\gamma_k}, \quad \text{where} \quad \gamma_k := \frac{\psi(k)}{\sqrt{k}}. \quad (12)$$

Note that  $\mathcal{K}_n$  is never negative. By the upper bound in (5), as  $k \rightarrow \infty$  we have

$$\gamma_k \leq \frac{\psi^U(k)}{\sqrt{k}} = \sqrt{\frac{2 \ln_2 k + 4 \ln_3 k}{k}} \rightarrow 0. \quad (13)$$

We show that  $\limsup_n \mathcal{K}_n = \infty$  if a path  $\omega \in \Omega_{<\infty}$  satisfies  $S_n \geq A_n \psi(A_n^2)$  i.o. We bound  $Z\mathcal{K}_n$  as

$$Z\mathcal{K}_n \geq \sum_{k=\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor}^{\lfloor A_n^2 \rfloor} p_k \mathcal{K}_n^{\gamma_k}. \quad (14)$$

At this point we check that all accounts on the right-hand side of (14) are open for sufficiently large  $n$  and the lower bound in (9) can be applied to each term of (14) for  $\omega \in \Omega_{<\infty}$ . We have the following two lemmas.

**Lemma 3.1.** *Let  $\omega \in \Omega_{<\infty}$ . Let  $C = C(\omega)$  in (10). For sufficiently large  $n$*

$$\bar{c}_n = \max_{1 \leq i \leq n} c_i < (1 + \delta) C \frac{A_n}{\psi(A_n^2)^3}. \quad (15)$$

*Proof.* Note that the first  $n_1(\omega)$   $c$ 's i.e.,  $c_1, \dots, c_{n_1(\omega)}$ , do not matter since  $\lim_{n \rightarrow \infty} A_n/\psi(A_n^2)^3 = \infty$ . For  $l > n_1(\omega)$ , by (10) we have

$$c_l \leq C \frac{A_l}{\psi(A_l^2)^3} \leq CA_l.$$

Hence  $c_l$  such that  $A_l \leq A_n/\psi(A_n^2)^3$  do not matter in  $\bar{c}_n$ .

For  $c_l$  such that  $A_l > A_n/\psi(A_n^2)^3$  we have

$$c_l \leq C \frac{A_l}{\psi(A_n^2/\psi(A_n^2)^6)^3} \leq C \frac{A_n}{\psi(A_n^2/\psi(A_n^2)^6)^3} = C \frac{A_n}{\psi(A_n^2)^3} \frac{\psi(A_n^2)^3}{\psi(A_n^2/\psi(A_n^2)^6)^3}.$$

But by (5), both  $\psi(A_n^2)$  and  $\psi(A_n^2/\psi(A_n^2)^6)$  are of the order  $\sqrt{2 \ln_2 A_n^2}(1+o(1))$  and  $\psi(A_n^2)/\psi(A_n^2/\psi(A_n^2)^6) \rightarrow 1$  as  $n \rightarrow \infty$ . Hence (15) holds.  $\square$

**Lemma 3.2.** *Let  $\omega \in \Omega_{<\infty}$ . For sufficiently large  $n$ ,  $\sigma_{\gamma_k} > n$  for all  $k = \lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor, \dots, \lfloor A_n^2 \rfloor$ .*

*Proof.* By the monotonicity of  $\psi$ , we have  $\gamma_k \leq \psi(A_n^2)/\sqrt{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor}$  for  $k = \lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor, \dots, \lfloor A_n^2 \rfloor$ . Then by the monotonicity of  $\sigma_\gamma$ , it suffices to show

$$\frac{\psi(A_n^2)}{\sqrt{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor}} \bar{c}_n \leq \delta$$

for sufficiently large  $n$ . By (15), the left-hand side is bounded from above by

$$\frac{\psi(A_n^2)}{\sqrt{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor}} \times (1 + \delta) C \frac{A_n}{\psi(A_n^2)^3} = (1 + \delta) C \frac{A_n}{\sqrt{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor}} \frac{1}{\psi(A_n^2)^2}.$$

But this converges to 0 as  $n \rightarrow \infty$ .  $\square$

By Lemma 3.2 and the lower bound in (9), for sufficiently large  $n$ , we have

$$\mathcal{K}_n^{\gamma_k} \geq e^{-\gamma_k^3 A_n^2 \bar{c}_n} e^{\gamma_k S_n - \gamma_k^2 A_n^2 / 2}, \quad k = \lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor, \dots, \lfloor A_n^2 \rfloor$$

and  $Z\mathcal{K}_n$  can be evaluated from below as

$$\begin{aligned} Z\mathcal{K}_n &\geq Z \sum_{k=\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor}^{\lfloor A_n^2 \rfloor} p_k \exp(\gamma_k S_n - \frac{\gamma_k^2 A_n^2}{2} - \gamma_k^3 A_n^2 \bar{c}_n) \\ &= \sum_{k=\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor}^{\lfloor A_n^2 \rfloor} a_k \frac{\psi(k)}{k} \exp(-\frac{\psi(k)^2}{2} + \gamma_k S_n - \frac{\gamma_k^2 A_n^2}{2} - \gamma_k^3 A_n^2 \bar{c}_n) \end{aligned}$$

Now we assume that  $S_n \geq A_n \psi(A_n^2)$  i.o. for the path  $\omega \in \Omega_{<\infty}$ . Then for sufficiently large  $n$  such that  $S_n \geq A_n \psi(A_n^2)$ ,  $\psi(A_n^2) / (\psi(A_n^2) - 1) \leq 1 + \delta$  and  $A_n / \left( \lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor \right)^{1/2} \leq 1 + \delta$ , we evaluate the exponent part by (9) as

$$\begin{aligned} -\frac{\psi(k)^2}{2} + \gamma_k S_n - \frac{\gamma_k^2 A_n^2}{2} &\geq -\frac{\psi(k)^2}{2} + A_n \psi(A_n^2) \frac{\psi(k)}{\sqrt{k}} - \frac{\psi(k)^2 A_n^2}{k} \\ &= \psi(k) \left( -\frac{1}{2} \left( 1 + \frac{A_n^2}{k} \right) \psi(k) + \sqrt{\frac{A_n^2}{k}} \psi(A_n^2) \right) \\ &\geq -\frac{\psi(A_n^2)^2}{2} \left( \sqrt{\frac{A_n^2}{k}} - 1 \right)^2 \geq -\frac{\psi(A_n^2)^2}{2} \left( \frac{A_n^2}{k} - 1 \right)^2 \\ &\geq -\frac{1}{2} \left( \frac{\psi(A_n^2)}{\psi(A_n^2) - 1} \right)^2 \geq -\frac{1}{2} - 2\delta \end{aligned}$$

and by Lemma 3.1

$$\begin{aligned} \gamma_k^3 A_n^2 \bar{c}_n &\leq \frac{\psi(A_n^2)^3}{(\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor)^{3/2}} A_n^2 (1 + \delta) C \frac{A_n}{\psi(A_n^2)^3} \\ &\leq (1 + \delta) C \left( \frac{A_n}{(\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor)^{1/2}} \right)^3 \\ &\leq C(1 + \delta)^4. \end{aligned} \tag{16}$$

For sufficiently large  $n$ , we have

$$\psi(A_n^2) \leq \psi^U(A_n^2) < \psi^U(2k) = \sqrt{2 \ln_2 2k + 4 \ln_3 2k} < 2 \sqrt{2 \ln_2 k + 3 \ln_2 k} = 2\psi^L(k) \leq 2\psi(k).$$

Thus by (16),

$$\begin{aligned} Z\mathcal{K}_n &\geq \sum_{k=\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor}^{\lfloor A_n^2 \rfloor} a_k \frac{\psi(k)}{k} \exp\left(-\frac{1}{2} - 2\delta - C(1 + \delta)^4\right) \\ &\geq a_{\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor} \frac{\psi(A_n^2)}{2A_n^2} \sum_{k=\lfloor A_n^2 - A_n^2 / \psi(A_n^2) \rfloor}^{\lfloor A_n^2 \rfloor} \exp\left(-\frac{1}{2} - 2\delta - C(1 + \delta)^4\right) \end{aligned}$$

$$\begin{aligned}
&\geq a_{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor} \frac{\psi(A_n^2)}{2A_n^2} \left( \frac{A_n^2}{\psi(A_n^2)} - 1 \right) \exp\left(-\frac{1}{2} - 2\delta - C(1+\delta)^4\right) \\
&= a_{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor} \left( \frac{1}{2} - \frac{\psi(A_n^2)}{2A_n^2} \right) \exp\left(-\frac{1}{2} - 2\delta - C(1+\delta)^4\right).
\end{aligned}$$

Since  $a_{\lfloor A_n^2 - A_n^2/\psi(A_n^2) \rfloor} \rightarrow \infty$  as  $n \rightarrow \infty$ , we have shown

$$\omega \in \Omega_{<\infty}, S_n \geq A_n \psi(A_n^2) \text{ i.o.} \Rightarrow \limsup_{n \rightarrow \infty} \mathcal{K}_n = \infty.$$

## 4 Sharpness

We prove the sharpness in (7) of Theorem 2.1. As in Section 4.2 of [17] and in [13], in order to prove the sharpness, it suffices to show the following proposition.

**Proposition 4.1.** *Consider SPUFG. Let  $\psi$  be a positive non-decreasing continuous function. If  $I(\psi) = \infty$ , then for each  $C > 0$ , Skeptic can force*

$$A_n^2 \rightarrow \infty, \limsup_n c_n \frac{\psi(A_n^2)^3}{A_n} \leq C \Rightarrow S_n \geq A_n \psi(A_n^2) \text{ i.o.} \quad (17)$$

Once we prove this proposition, we can take the mixture over  $C = 1, 2, \dots$ . Then the sharpness follows, because for each  $\omega \in \Omega_{<\infty}$ , there exists  $C(\omega)$  satisfying (10). We denote

$$\begin{aligned}
\Omega_C &:= \left\{ \omega \in \Omega \mid A_n^2 \rightarrow \infty, \limsup_n c_n \frac{\psi(A_n^2)^3}{A_n} < (1-\delta)C \right\}, \\
\Omega_0 &:= \left\{ \omega \in \Omega \mid \lim_{n \rightarrow \infty} A_n^2 < \infty \right\}, \\
\Omega_{=\infty} &:= \left\{ \omega \in \Omega \mid A_n^2 \rightarrow \infty, \limsup_n c_n \frac{\psi(A_n^2)^3}{A_n} = \infty \right\}.
\end{aligned}$$

We divide our proof of Proposition 4.1 into several subsections. For notational simplicity we use the initial capital of  $\alpha = 1 - 2/e = (e-2)/e$  in this section. In Sections 4.1 and 4.2 we only consider  $\gamma$  and  $n$  with  $n < \sigma_\gamma$ . As in Lemma 3.2 for the validity, this condition will be satisfied for sufficiently small  $\gamma$  and relevant  $n$ .

### 4.1 Uniform mixture of constant-proportion betting strategies

We consider a continuous uniform mixture of constant-proportion strategies with the betting proportion  $u\gamma$ ,  $2/e \leq u \leq 1$ . This is a Bayesian strategy, a similar one to which has been considered in [12].

Define

$$\mathcal{L}_n^\gamma := \int_{2/e}^1 \prod_{i=1}^{\min(n, \sigma_\gamma - 1)} (1 + u\gamma x_i) du, \quad \mathcal{L}_0^\gamma = \alpha = 1 - e/2.$$

At round  $n < \sigma_\gamma$  this strategy bets  $M_n = \int_{2/e}^1 u\gamma \prod_{i=1}^{n-1} (1 + u\gamma x_i) du$ . Then by induction on  $n < \sigma_\gamma$  the capital process is indeed written as

$$\mathcal{L}_n^\gamma = \mathcal{L}_{n-1}^\gamma + M_n x_n = \int_{2/e}^1 \prod_{i=1}^{n-1} (1 + u\gamma x_i) du + x_n \int_{2/e}^1 u\gamma \prod_{i=1}^{n-1} (1 + u\gamma x_i) du$$

$$= \int_{2/e}^1 \prod_{i=1}^n (1 + u\gamma x_i) du.$$

Applying (9), we have

$$e^{-\gamma^3 A_n^2 \bar{c}_n} \int_{2/e}^1 e^{u\gamma S_n - u^2 \gamma^2 A_n^2 / 2} du \leq \mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} \int_{2/e}^1 e^{u\gamma S_n - u^2 \gamma^2 A_n^2 / 2} du,$$

for  $n < \sigma_\gamma$ . We further bound the integral in the following lemma.

**Lemma 4.2.** For  $n < \sigma_\gamma$ ,

$$\mathcal{L}_n^\gamma \leq \begin{cases} e^{\gamma^3 A_n^2 \bar{c}_n} e^{2\gamma(S_n/e - \gamma A_n^2/e^2)} & \text{if } S_n \leq 2\gamma A_n^2/e, \\ e^{\gamma^3 A_n^2 \bar{c}_n} \min \left\{ e^{S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n/2} \right\} & \text{if } 2\gamma A_n^2/e < S_n < \gamma A_n^2, \\ e^{\gamma^3 A_n^2 \bar{c}_n} \min \left\{ e^{S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n - \gamma^2 A_n^2/2} \right\} & \text{if } S_n \geq \gamma A_n^2. \end{cases} \quad (18)$$

$$\mathcal{L}_n^\gamma \leq \begin{cases} e^{\gamma^3 A_n^2 \bar{c}_n} \min \left\{ e^{S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n/2} \right\} & \text{if } 2\gamma A_n^2/e < S_n < \gamma A_n^2, \end{cases} \quad (19)$$

$$\mathcal{L}_n^\gamma \leq \begin{cases} e^{\gamma^3 A_n^2 \bar{c}_n} \min \left\{ e^{S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n - \gamma^2 A_n^2/2} \right\} & \text{if } S_n \geq \gamma A_n^2. \end{cases} \quad (20)$$

*Proof.* Completing the square we have

$$-\frac{1}{2}u^2\gamma^2 A_n^2 + u\gamma S_n = -\frac{\gamma^2 A_n^2}{2} \left( u - \frac{S_n}{\gamma A_n^2} \right)^2 + \frac{S_n^2}{2A_n^2}.$$

Hence by the change of variables

$$v = \gamma A_n \left( u - \frac{S_n}{\gamma A_n^2} \right), \quad du = \frac{dv}{\gamma A_n},$$

we obtain

$$\begin{aligned} \int_{2/e}^1 e^{u\gamma S_n - u^2 \gamma^2 A_n^2 / 2} du &= e^{S_n^2/(2A_n^2)} \int_{2/e}^1 \exp \left( -\frac{\gamma^2 A_n^2}{2} \left( u - \frac{S_n}{\gamma A_n^2} \right)^2 \right) du \\ &= e^{S_n^2/(2A_n^2)} \frac{1}{\gamma A_n} \int_{2\gamma A_n/e - S_n/A_n}^{\gamma A_n - S_n/A_n} e^{-v^2/2} dv. \end{aligned}$$

Then for all cases we can bound  $\mathcal{L}_n^\gamma$  from above as

$$\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n + S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}. \quad (21)$$

Without change of variables, we can also bound the integral  $\int_{2/e}^1 g(u) du$ ,  $g(u) := e^{u\gamma S_n - u^2 \gamma^2 A_n^2 / 2}$ , directly as

$$\int_{2/e}^1 g(u) du \leq \max_{2/e \leq u \leq 1} g(u).$$

Note that

$$g(2/e) = e^{2\gamma(S_n/e - \gamma A_n^2/e^2)}, \quad g(1) = e^{\gamma S_n - \gamma^2 A_n^2/2}. \quad (22)$$

We now consider the following three cases.



**Case 1**  $S_n \leq 2\gamma A_n^2/e$ . In this case  $S_n/(\gamma A_n^2) \leq 2/e$  and by the unimodality of  $g(u)$  we have  $\max_{2/e \leq u \leq 1} g(u) = g(2/e)$ . Hence (18) follows from (22).

**Case 2**  $2\gamma A_n^2/e < S_n < \gamma A_n^2$ . In this case  $\max_{2/e \leq u \leq 1} g(u) = g(S_n/(\gamma A_n^2)) = e^{S_n^2/(2A_n^2)}$  and  $\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{S_n^2/(2A_n^2)}$ . Furthermore in this case  $S_n^2 < \gamma A_n^2 S_n$  implies  $S_n^2/(2A_n^2) < \gamma S_n/2$  and we also have

$$\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n/2}. \quad (23)$$

By (21) and (23), we have (19).

**Case 3**  $S_n \geq \gamma A_n^2$ . Then  $S_n/(\gamma A_n^2) \geq 1$  and  $\max_{2/e \leq u \leq 1} g(u) = g(1)$ . Hence

$$\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2/2}. \quad (24)$$

By (21) and (24), we have (20). □

## 4.2 Buying a process and selling a process

Next we consider the following capital process.

$$Q_n^\gamma := 2\mathcal{L}_n^\gamma - \mathcal{K}_n^{\gamma e}. \quad (25)$$

This capital process consists of buying two units of  $\mathcal{L}_n^\gamma$  and selling one unit of  $\mathcal{K}_n^{\gamma e}$ . This combination of selling and buying is essential in the game-theoretic proof of LIL in Chapter 5 of [17] and [14]. However, unlike Chapter 5 of [17] and [14], where a combination of *three* capital processes is used, we only combine *two* capital processes.

We want to bound  $Q_n^\gamma$  from above.

**Lemma 4.3.** *Let*

$$C_1 := 2e^{\gamma^3 A_n^2 \bar{c}_n} \exp\left(\frac{(2e-1)((1+e^3)\gamma^3 A_n^2 \bar{c}_n + \ln 2)}{(e-1)^2}\right). \quad (26)$$

*Then for  $n < \sigma_{\gamma e}$ ,*

$$Q_n^\gamma \leq \begin{cases} C_1 & \text{if } S_n \leq \gamma A_n^2/e, \end{cases} \quad (27)$$

$$Q_n^\gamma \leq \begin{cases} 2e^{\gamma^3 A_n^2 \bar{c}_n} \min\left\{e^{S_n^2/(2A_n^2)} \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n}\right\} & \text{if } \gamma A_n^2/e < S_n < e\gamma A_n^2, \end{cases} \quad (28)$$

$$Q_n^\gamma \leq \begin{cases} C_1 & \text{if } S_n \geq e\gamma A_n^2. \end{cases} \quad (29)$$

**Remark 4.4.** In this lemma,  $C_1$  depends on  $\bar{c}_n$ ,  $\gamma$  and  $A_n$  through  $\gamma^3 A_n^2 \bar{c}_n$ . However from Section 4.5 on, we evaluate  $\gamma^3 A_n^2 \bar{c}_n$  from above by a constant. Hence,  $C_1$  can be also taken to be a constant (cf. (50)) not depending on  $\gamma$  and  $A_n$ . Also note that the interval for  $S_n$  in (28) is larger than the interval in (19).

*Proof.* We bound  $Q_n^\gamma = 2\mathcal{L}_n^\gamma - \mathcal{K}_n^{\gamma e}$  from above in the following three cases:

$$(i) S_n \leq \gamma A_n^2/e, \quad (ii) \gamma A_n^2/e < S_n < e\gamma A_n^2, \quad (iii) S_n \geq e\gamma A_n^2,$$

**Case (i)** In this case  $S_n/e - \gamma A_n^2/e^2 \leq 0$ . Hence (27) follows from (18) and  $Q_n^\gamma \leq 2\mathcal{L}_n^\gamma$ .

**Case (ii)** We again use  $Q_n^\gamma \leq 2\mathcal{L}_n^\gamma$ . If  $\gamma A_n^2/e < S_n \leq 2\gamma A_n^2/e$ , then

$$\frac{S_n}{e} - \frac{\gamma A_n^2}{e^2} \leq \frac{\gamma A_n^2}{e^2} \leq \frac{S_n}{e}$$

and  $\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{2\gamma S_n/e} \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n}$  from (18). Otherwise (28) follows from (19) and (20).

**Case (iii)** Since  $S_n \geq eA_n^2\gamma > A_n^2\gamma$ , by (24) we have  $\mathcal{L}_n^\gamma \leq e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2/2}$  and

$$\begin{aligned} Q_n^\gamma &\leq 2\mathcal{L}_n^\gamma - \mathcal{K}_n^{\gamma e} \leq 2e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2/2} - e^{-\gamma^3 e^3 A_n^2 \bar{c}_n} e^{\gamma e S_n - \gamma^2 e^2 A_n^2/2} \\ &= 2e^{\gamma^3 A_n^2 \bar{c}_n} e^{\gamma S_n - \gamma^2 A_n^2/2} \left( 1 - \frac{1}{2} e^{-(1+e^3)\gamma^3 A_n^2 \bar{c}_n} e^{\gamma(e-1)S_n - (e^2-1)\gamma^2 A_n^2/2} \right). \end{aligned}$$

Hence if the right-hand side is non-positive we have  $Q_n^\gamma \leq 0$ :

$$\begin{aligned} S_n &\geq eA_n^2\gamma \quad \text{and} \quad -(1+e^3)\gamma^3 A_n^2 \bar{c}_n - \ln 2 + \gamma(e-1)S_n - \frac{1}{2}(e^2-1)\gamma^2 A_n^2 \geq 0 \\ &\Rightarrow Q_n^\gamma \leq 0. \end{aligned} \tag{30}$$

Otherwise, write  $B_n := (1+e^3)\gamma^3 A_n^2 \bar{c}_n + \ln 2$  and consider the case

$$\gamma(e-1)S_n - \frac{1}{2}(e^2-1)\gamma^2 A_n^2 \leq B_n.$$

Dividing this by  $e-1$  and also considering  $S_n \geq eA_n^2\gamma$ , we have

$$\gamma S_n - \frac{1}{2}(e+1)\gamma^2 A_n^2 \leq \frac{B_n}{e-1}, \tag{31}$$

$$-S_n + eA_n^2\gamma \leq 0. \tag{32}$$

$\gamma \times (32) + (31)$  gives

$$\frac{1}{2}(e-1)\gamma^2 A_n^2 \leq \frac{B_n}{e-1} \quad \text{or} \quad \frac{1}{2}\gamma^2 A_n^2 \leq \frac{B_n}{(e-1)^2}.$$

Then by (31)

$$\gamma S_n - \frac{1}{2}\gamma^2 A_n^2 \leq \frac{B_n}{e-1} + \frac{e}{2}\gamma^2 A_n^2 \leq \frac{B_n}{e-1} + \frac{eB_n}{(e-1)^2} = \frac{(2e-1)B_n}{(e-1)^2}.$$

Hence just using  $Q_n^\gamma \leq 2\mathcal{L}_n^\gamma$  and (24) in this case, we obtain

$$Q_n^\gamma \leq 2e^{\gamma^3 A_n^2 \bar{c}_n} \exp\left(\frac{(2e-1)((1+e^3)\gamma^3 A_n^2 \bar{c}_n + \ln 2)}{(e-1)^2}\right) = C_1. \tag{33}$$

This also covers (30) and we have (33) for the whole case (iii).

□

### 4.3 Change of time scale and dividing the rounds into cycles

For proving the sharpness we consider the change of time scale from  $\lambda$  to  $k$ :

$$\lambda = e^{5k \ln k} = k^{5k}.$$

By taking the derivative of  $\ln \lambda = 5k \ln k$ , we have  $d\lambda/\lambda = 5(\ln k + 1)dk$ . Since  $\ln k$  is dominant in  $(\ln k + 1)$ , the integrability condition is written as

$$\int_1^\infty \psi(\lambda) e^{-\psi(\lambda)^2/2} \frac{d\lambda}{\lambda} = \infty \Leftrightarrow \int_1^\infty (\ln k) \psi(e^{5k \ln k}) e^{-\psi(e^{5k \ln k})^2/2} dk = \infty.$$

Let  $f(x) := \psi(e^{5x \ln x}) e^{-\psi(e^{5x \ln x})^2/2}$ . Since  $xe^{-x^2/2}$  is decreasing for  $x \geq 1$ , the function  $f(x)$  is decreasing for  $x$  such that  $\psi(e^{5x \ln x}) \geq 1$ . Thus, for sufficiently large  $k$  and  $x$  such that  $k \leq x \leq k+1$ , we have

$$\frac{1}{2} \ln(k+1) f(k+1) \leq \ln k f(x+1) \leq \ln x f(x) \leq \ln(k+1) f(x) \leq 2 \ln k f(k).$$

Hence, we have

$$\int_1^\infty (\ln k) \psi(e^{5k \ln k}) e^{-\psi(e^{5k \ln k})^2/2} dk = \infty \Leftrightarrow \sum_{k=1}^\infty (\ln k) \psi(e^{5k \ln k}) e^{-\psi(e^{5k \ln k})^2/2} = \infty.$$

Then, it suffices to show (17) if  $\sum_{k=1}^\infty (\ln k) \psi(e^{5k \ln k}) e^{-\psi(e^{5k \ln k})^2/2} = \infty$ .

As in Chapter 5 of [17] and [14], we divide the time axis into “cycles”. However, unlike in Chapter 5 of [17] and [14], our cycles are based on stopping times. Let

$$n_k := k^{5k}, \quad k = 1, 2, \dots, \quad (34)$$

and define a family of stopping times

$$\tau_k := \min \{n \mid A_n^2 \geq n_k\}. \quad (35)$$

We define the  $k$ -th cycle by  $[\tau_k, \tau_{k+1}]$ ,  $k \geq 1$ . Note that  $\tau_k$  is finite for all  $k$  if and only if  $A_n^2 \rightarrow \infty$ . Betting strategy for the  $k$ -th cycle is based on the following betting proportion:

$$\gamma_k := \frac{\psi(n_{k+1})}{\sqrt{n_{k+1}}} k^2. \quad (36)$$

Note that  $\gamma_k$  in (36) is slightly different from (12).

For the rest of this section, we check the growth of various quantities along the cycles. Let  $\omega \in \Omega_C$ . For sufficiently large  $n$ ,

$$|x_n| \leq c_n \leq C \frac{A_n}{\psi(A_n^2)^3}. \quad (37)$$

Furthermore  $A_n^2 = A_{n-1}^2 + x_n^2$ . This allows us to bound  $x_n^2$  and  $A_n^2$  in terms of  $A_{n-1}^2$ . By squaring (37) we have

$$x_n^2 \leq C^2 \frac{A_{n-1}^2}{\psi(A_n^2)^6 - C^2} \quad (38)$$

and

$$A_n^2 = A_{n-1}^2 + x_n^2 \leq A_{n-1}^2 \left(1 + \frac{C^2}{\psi(A_n^2)^6 - C^2}\right) = A_{n-1}^2 \frac{\psi(A_n^2)^6}{\psi(A_n^2)^6 - C^2}. \quad (39)$$

Since  $\psi(A_n^2)^6/(\psi(A_n^2)^6 - C^2) \rightarrow 1$  as  $n \rightarrow \infty$ , we have

$$\lim_{n \rightarrow \infty} \frac{A_n^2}{A_{n-1}^2} = 1.$$

Note that  $A_{\tau_k-1}^2 < n_k \leq A_{\tau_k}^2$  by the definition of  $\tau_k$ . Hence for  $\omega \in \Omega_C$  we also have

$$\lim_{k \rightarrow \infty} \frac{A_{\tau_k}^2}{n_k} = 1. \quad (40)$$

The limits in the following lemma will be useful for our argument.

**Lemma 4.5.** *For  $\omega \in \Omega_C$*

$$\lim_{k \rightarrow \infty} \frac{\psi^U(n_k)}{\psi(n_{k+1})} = 1, \quad \lim_{k \rightarrow \infty} \frac{k^5 A_{\tau_k}^2}{n_{k+1}} = e^{-5}, \quad \lim_{k \rightarrow \infty} \gamma_k A_{\tau_k} \psi(n_{k+1}) = 0. \quad (41)$$

*Proof.* All of  $\psi^U(n_k)$ ,  $\psi^U(n_{k+1})$ ,  $\psi^L(n_k)$ ,  $\psi^L(n_{k+1})$ ,  $\psi(n_{k+1})$ ,  $\psi(n_{k+1}/k^4)$  are of the order

$$\sqrt{2 \ln \ln e^{5k \ln k}} (1 + o(1)) = \sqrt{2 \ln k} (1 + o(1)) \quad (42)$$

as  $k \rightarrow \infty$  and the first equality holds by (5). The second equality holds by (40) and

$$\lim_{k \rightarrow \infty} \frac{k^5 n_k}{n_{k+1}} = \lim_{k \rightarrow \infty} \frac{k^{5(k+1)}}{(k+1)^{5(k+1)}} = \lim_{k \rightarrow \infty} \left(1 - \frac{1}{k+1}\right)^{5(k+1)} = e^{-5}.$$

Then  $A_{\tau_k}^2/n_{k+1} = (1 + o(1))n_k/n_{k+1} = O(k^{-5})$  and the third equality holds by

$$\gamma_k A_{\tau_k} \psi(n_{k+1}) \leq \psi(n_{k+1})^2 k^2 ((1 + \delta)n_k/n_{k+1})^{1/2} \rightarrow 0 \quad (k \rightarrow \infty).$$

□

## 4.4 Stopping times for aborting and sequential freezing for each cycle

In (48) of the next section we will introduce another capital process  $\mathcal{M}_n^{\gamma_k, k}$ , which will be employed in each cycle. Here we introduce some stopping times for aborting the cycle and for sequential freezing of accounts in  $\mathcal{M}_n^{\gamma_k, k}$ .

We say that we *abort* the  $k$ -th cycle, when we freeze all accounts in the  $k$ -th cycle and wait for the  $(k+1)$ -st cycle. There are two cases for aborting the  $k$ -th cycle. The first case is when some  $c_n$  is too large for  $\omega \in \Omega_C$ . Define

$$\sigma_{k,C} := \min \{n \geq \tau_k \mid c_n \psi(A_{\tau_k}^2)^3 > (1 + \delta) C A_{n-1}\}. \quad (43)$$

We will abort the  $k$ -th cycle if  $\sigma_{k,C} < \tau_{k+1}$ . Note that for  $\omega \in \Omega_C$ , there exists  $k_1(\omega)$  such that

$$\sigma_{k,C} = \infty, \quad \text{for } k \geq k_1(\omega). \quad (44)$$

Another case is when  $S_n$  is too large. Define

$$\nu_k := \min \{n \geq \tau_k \mid A_n \psi(A_n^2) < S_n\}. \quad (45)$$

If  $\nu_k < \tau_{k+1}$ , then Skeptic is happy to abort the  $k$ -th cycle, because he wants to force  $S_n \geq A_n \psi(A_n^2)$  i.o. The above two stopping times will be used in the final construction of a dynamic strategy in Section 4.6.

For each cycle, we define another family of stopping times indexed by  $w = 1, \dots, \lceil \ln k \rceil$ , by

$$\tau_{k,w} := \min \left\{ n \mid A_n^2 \geq e^{2(w+2)} \frac{n_{k+1}}{k^4} \right\}. \quad (46)$$

for sequential freezing of accounts of  $\mathcal{M}_n^{\gamma_k, k}$  in (48). We have  $\tau_k \leq \tau_{k,w}$  for  $k \geq 1$  and  $w \geq 1$ , because

$$\frac{n_{k+1}}{k^4} = \frac{(k+1)^{5(k+1)}}{k^4} > k^{5k} = n_k.$$

**Lemma 4.6.** *Let  $\omega \in \Omega_C$ .  $\tau_{k, \lceil \ln k \rceil} \leq \tau_{k+1}$  for sufficiently large  $k$ .*

*Proof.* By  $A_{\tau_{k,w-1}}^2 \leq e^{2(w+2)} n_{k+1}/k^4$  and by (38), for sufficiently large  $k$  we have

$$x_{\tau_{k,w}}^2 \leq (1+\delta) C^2 \frac{A_{\tau_{k,w-1}}^2}{\psi(A_{\tau_k}^2)^6} \leq \frac{(1+\delta) C^2}{\psi(A_{\tau_k}^2)^6} \times \frac{e^{2(w+2)} n_{k+1}}{k^4}$$

and

$$A_{\tau_{k,w}}^2 \leq A_{\tau_{k,w-1}}^2 + x_{\tau_{k,w}}^2 \leq (1+\delta) e^{2(w+2)} \frac{n_{k+1}}{k^4}. \quad (47)$$

Then

$$A_{\tau_{k, \lceil \ln k \rceil}}^2 \leq (1+\delta) \left( e^{2(\ln k + 2)} \frac{n_{k+1}}{k^4} \right) = (1+\delta) e^4 \frac{n_{k+1}}{k^2} \leq n_{k+1} \leq A_{\tau_{k+1}}^2.$$

□

We also compare  $\tau_{k,w}$  to  $\sigma_{\gamma_k e^{-w+1}}$  defined in (8). This is needed for applying the bounds derived in previous sections to  $\mathcal{M}_n^{\gamma_k, k}$  in the next section.

**Lemma 4.7.** *Let  $\omega \in \Omega_C$ .  $\tau_{k,w} \leq \sigma_{\gamma_k e^{-w+1}}$  for sufficiently large  $k$ .*

*Proof.* By (47) and by Lemma 3.1, for sufficiently large  $k$

$$\gamma_k e^{-w+1} \bar{c}_{\tau_{k,w}} \leq \frac{\psi(n_{k+1})}{\sqrt{n_{k+1}}} k^2 e^{-w+1} \times (1+\delta)^2 C \frac{e^{w+2} \sqrt{n_{k+1}}}{k^2 \psi(A_{\tau_k}^2)^3} \leq (1+\delta)^2 C e^3 \frac{\psi(n_{k+1})}{\psi(A_{\tau_k}^2)^3} \leq \delta,$$

because  $\psi(n_{k+1})/\psi(A_{\tau_k}^2)^3 \rightarrow 0$  as  $k \rightarrow \infty$  by (42). □

## 4.5 Further discrete mixture of processes for each cycle with sequential freezing

We introduce another discrete mixture of capital process for the  $k$ -th cycle. Define

$$\mathcal{M}_n^{\gamma_k, k} := \frac{1}{\lceil \ln k \rceil} \sum_{w=1}^{\lceil \ln k \rceil} Q_{\min(n, \tau_{k,w})}^{\gamma_k e^{-w}} = \frac{1}{\lceil \ln k \rceil} \sum_{w=1}^{\lceil \ln k \rceil} (2\mathcal{L}_{\min(n, \tau_{k,w})}^{\gamma_k e^{-w}} - \mathcal{K}_{\min(n, \tau_{k,w})}^{\gamma_k e^{-w+1}}). \quad (48)$$

Note that the  $w$ -th account in the sum of  $\mathcal{M}_n^{\gamma_k, k}$  is frozen at the stopping time  $\tau_{k,w}$ . This is needed since the bound for  $c_n$  is growing even during the  $k$ -th cycle.

In order to bound  $\mathcal{M}_n^{\gamma_k, k}$ , we first bound  $C_1$  in (26) for each  $w$  in the sum of (48) by a constant independent of  $n$ . Note that we only need to consider  $n \leq \tau_{k,w}$  for the  $w$ -th account.

**Lemma 4.8.** Let  $\omega \in \Omega_C$ .  $(\gamma_k e^{-w})^3 A_n^2 \bar{c}_n$  and hence  $C_1$  are bounded from above by

$$(\gamma_k e^{-w})^3 A_n^2 \bar{c}_n \leq (1 + \delta)^5 C e^6, \quad (49)$$

$$C_1 \leq 2e^{(1+\delta)^5 C e^6} \exp\left(\frac{(2e-1)((1+\delta)^5 C e^6(1+e^3) + \ln 2)}{(e-1)^2}\right) =: \bar{C}_1, \quad (50)$$

for sufficiently large  $k$ .

*Proof.* By (42), for sufficiently large  $k$

$$\frac{\psi(n_{k+1})}{\psi(A_{\tau_{k,w}}^2)} \leq \frac{\psi(n_{k+1})}{\psi(n_k)} \leq 1 + \delta. \quad (51)$$

Thus

$$\begin{aligned} \gamma_k^3 e^{-3w} A_{\min(n, \tau_{k,w})}^2 \bar{c}_{\min(n, \tau_{k,w})} &\leq \gamma_k^3 e^{-3w} \times A_{\tau_{k,w}}^2 \times \bar{c}_{\min(n, \tau_{k,w})} \\ &\leq \frac{\psi(n_{k+1})^3}{n_{k+1}^{3/2}} k^6 e^{-3w} \times A_{\tau_{k,w}}^2 \times (1 + \delta) C \frac{A_{\tau_{k,w}}}{\psi(A_{\tau_k}^2)^3} \\ &\leq (1 + \delta) C \frac{\psi(n_{k+1})^3}{\psi(A_{\tau_k}^2)^3} k^6 e^{-3w} \frac{A_{\tau_{k,w}}^3}{n_{k+1}^{3/2}} \leq (1 + \delta)^5 C e^6. \end{aligned}$$

□

**Lemma 4.9.** Let  $\omega \in \Omega_C$ . For sufficiently large  $k$ ,

$$\mathcal{M}_n^{\gamma_k, k} \leq \bar{C}_1 + \frac{2}{\lceil \ln k \rceil} e^{(1+\delta)^5 C e^6} \max_{\gamma \in [\gamma_k/k, \gamma_k]} \left( \min\{e^{S_n^2/(2n)}, \frac{\sqrt{2\pi}}{\gamma A_n}, e^{\gamma S_n}\} \right), \quad n \in [\tau_k, \tau_{k+1}], \quad (52)$$

where  $\bar{C}_1$  is given by the right-hand side of (50).

*Proof.* We have  $|\gamma_k e^{-w} \bar{c}_{\min(n, \tau_{k,w})}| \leq |\gamma_k e^{-w+1} \bar{c}_{\min(n, \tau_{k,w})}| \leq \delta$  by Lemma 4.7. Then we can complete the proof of (52) by Lemma 4.3 and Lemma 4.7 because the length of the interval

$$\left\{ w \mid \frac{S_n}{ne} < \gamma e^{-w} < \frac{S_n e}{n} \right\}$$

is equal to 2.

□

As in Chapter 5 of Shafer and Vovk [17], we use  $\mathcal{M}_n^{\gamma_k, k}$  in the following form.

$$\mathcal{N}_n^{\gamma_k, D} := \alpha + \frac{1}{D} \lceil \ln k \rceil \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} (\alpha - \mathcal{M}_{n-\tau_k}^{\gamma_k, k}), \quad \alpha = 1 - \frac{2}{e}, \quad D = \frac{24 \sqrt{2\pi} e^{(1+\delta)^5 e^6 C} + 4\bar{C}_1}{\alpha}. \quad (53)$$

Here we give a specific value of  $D$  for definiteness, but from the proof below it will be clear that any sufficiently large  $D$  can be used. Since the strategy for  $\mathcal{M}_{n-\tau_k}^{\gamma_k, k}$  is applied only to  $x_n$ 's in the cycle,  $\alpha = \mathcal{N}_{\tau_k}^{\gamma_k, D} = \mathcal{M}_0^{\gamma_k}$ . Concerning  $\mathcal{N}_n^{\gamma_k, D}$  we prove the following two propositions.

**Proposition 4.10.** Let  $\omega \in \Omega_C$ . Suppose that

$$-A_n \psi^U(A_n^2) \leq S_n \leq A_n \psi(A_n^2), \quad \forall n \in [\tau_k, \tau_{k+1}]. \quad (54)$$

and  $\tau_{k+1} < \sigma_{k,C}$ . Then for sufficiently large  $k$

$$\mathcal{N}_n^{\gamma_k, D} \geq \frac{\alpha}{2}, \quad \forall n \in [\tau_k, \tau_{k+1}], \quad (55)$$

and

$$\mathcal{N}_{\tau_{k+1}}^{\gamma_k, D} \geq \alpha \left( 1 + \frac{1-\delta}{D} \lceil \ln k \rceil \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} \right). \quad (56)$$

*Proof.* In our proof we denote  $t = n - \tau_k$ ,  $S_t = S_n - S_{\tau_k}$  and  $A_t^2 = A_n^2 - A_{\tau_k}^2$  for  $n > \tau_k$ . For proving (55), we use (52) for  $S_t$ . We bound  $\mathcal{M}_t^{\gamma_k, k}$  from above. By the term  $\frac{2}{\lceil \ln k \rceil}$  on the right-hand side of (52), it suffices to show

$$\begin{aligned} S_t &\leq A_{\tau_k} \psi^U(A_{\tau_k}^2) + \sqrt{A_{\tau_k}^2 + A_t^2} \psi(A_{\tau_k}^2 + A_t^2) \\ &\Rightarrow \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} 2e^{(1+\delta)^5 e^6 C} \min\{e^{S_t^2/(2A_t^2)} \frac{\sqrt{2\pi}}{\gamma A_t}, e^{\gamma S_t}\} \leq \frac{D\alpha}{4}, \quad \forall \gamma \in [\gamma_k/k, \gamma_k], \quad \forall t \in [0, \tau_{k+1} - \tau_k] \end{aligned}$$

for sufficient large  $k$ . Let

$$c_1 = \frac{9}{(1+2\delta)^2} \quad \text{s.t.} \quad \frac{1}{2} - \frac{1}{\sqrt{c_1}} - \delta > 0. \quad (57)$$

We distinguish two cases:

$$(a) A_t^2 \leq \frac{\psi(n_{k+1})^2}{c_1 \gamma^2}, \quad (b) \frac{\psi(n_{k+1})^2}{c_1 \gamma^2} < A_t^2 \leq A_{\tau_{k+1}}^2 - A_{\tau_k}^2.$$

For case (a),  $A_{\tau_k} \psi^U(A_{\tau_k}^2) \leq (1+\delta) A_{\tau_k} \psi(n_{k+1})$  by the first equality in Lemma 4.5 for sufficiently large  $k$ . Also  $\psi(A_{\tau_k}^2 + A_t^2) \leq \psi(n_{k+1})$ . Hence in this case

$$\gamma S_t \leq \left( (1+\delta) \gamma A_{\tau_k} + \sqrt{\gamma^2 A_{\tau_k}^2 + \psi(n_{k+1})^2/c_1} \right) \psi(n_{k+1}).$$

Then for  $\gamma \leq \gamma_k$  by the third equality in Lemma 4.5

$$\gamma S_t \leq \left( (1+\delta) \gamma_k A_{\tau_k} + \sqrt{\gamma_k^2 A_{\tau_k}^2 + \psi(n_{k+1})^2/c_1} \right) \psi(n_{k+1}) = \psi(n_{k+1})^2 \left( \frac{1}{\sqrt{c_1}} + \delta \right) \quad (58)$$

for sufficiently large  $k$ . Since

$$\psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} 2e^{(1+\delta)^5 e^6 C} e^{\gamma S_t} \leq \psi(n_{k+1}) \exp\left(-\psi(n_{k+1})^2 \left(\frac{1}{2} - \frac{1}{\sqrt{c_1}} - \delta\right)\right) 2e^{(1+\delta)^5 e^6 C} \rightarrow 0 \quad (k \rightarrow \infty),$$

we have  $\mathcal{N}_n^{\gamma_k, D} \geq \alpha/2$  uniformly in  $\gamma \in [\gamma_k/k, \gamma_k]$ .

For case (b),  $\psi(n_{k+1})/\sqrt{c_1} < \gamma A_t$  and  $S_t \leq \left( (1+\delta) A_{\tau_k} + \sqrt{A_{\tau_k}^2 + A_t^2} \right) \psi(n_{k+1})$ . Hence

$$\psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} \times 2e^{(1+\delta)^5 e^6 C} e^{S_t^2/(2A_t^2)} \frac{\sqrt{2\pi}}{\gamma A_t}$$

$$\begin{aligned}
&\leq \psi(n_{k+1})e^{-\psi(n_{k+1})^2/2} \times \frac{2e^{(1+\delta)^5 e^6 C} \sqrt{2\pi} \sqrt{c_1}}{\psi(n_{k+1})} \exp\left(\frac{\left((1+\delta)A_{\tau_k} + \sqrt{A_{\tau_k}^2 + A_t^2}\right)^2}{2A_t^2} \psi(n_{k+1})^2\right) \\
&= 2e^{(1+\delta)^5 e^6 C} \sqrt{2\pi} \sqrt{c_1} \exp\left(\frac{(1+(1+\delta)^2)A_{\tau_k}^2 + 2(1+\delta)A_{\tau_k} \sqrt{A_{\tau_k}^2 + A_t^2}}{2A_t^2} \psi(n_{k+1})^2\right). \tag{59}
\end{aligned}$$

For  $\gamma \leq \gamma_k$ ,

$$\frac{\psi(n_{k+1})^2}{c_1 \gamma^2} < A_t^2 \Rightarrow \frac{A_{\tau_k}^2}{A_t^2} \psi(n_{k+1})^2 < c_1 \gamma^2 A_{\tau_k}^2 \leq c_1 \gamma_k^2 A_{\tau_k}^2 = c_1 \frac{A_{\tau_k}^2}{n_{k+1}} k^4 \psi(n_{k+1})^2 = O(k^{-1} \ln k).$$

Hence  $\psi(n_{k+1})^2 A_{\tau_k}^2 / A_t^2 \rightarrow 0$  as  $k \rightarrow \infty$ . Similarly  $\psi(n_{k+1})^2 A_{\tau_k} / A_t \rightarrow 0$  as  $k \rightarrow \infty$ , because  $\psi(n_{k+1})^2 A_{\tau_k} / A_t = O(k^{-1/2} (\ln k)^{3/2})$ . Therefore the right-hand side of (59) is bounded from above by  $2e^{(1+\delta)^5 e^6 C} \sqrt{2\pi} \sqrt{c_1} (1+\delta)$  for sufficiently large  $k$  and

$$\psi(n_{k+1})e^{-\psi(n_{k+1})^2/2} \times 2e^{(1+\delta)^5 e^6 C} e^{S_t^2/(2A_t^2)} \frac{\sqrt{2\pi}}{\gamma A_t} \leq \frac{D\alpha}{4},$$

with the choice of  $D$  in (53) and  $c_1$  in (57). This proves (55).

Now we prove (56). We focus on the  $w$ -th account when  $n \geq \tau_{k,w}$ . Recall that in this proof we have been denoting  $A_t^2 = A_n^2 - A_{\tau_k}^2$ . Similarly we denote  $A_{\tau_{k,w}}^2$  instead of  $A_{\tau_{k,w}}^2 - A_{\tau_k}^2$ . Thus

$$e^{2(w+2)} \frac{n_{k+1}}{k^4} - A_{\tau_k}^2 \leq A_{\tau_{k,w}}^2. \tag{60}$$

We will show that  $\limsup_{k \rightarrow \infty} \mathcal{M}_{\tau_{k+1}-\tau_k}^{\gamma_k, k} \leq 0$ , if

$$S_{\tau_{k,w}} \leq A_{\tau_k} \psi(A_{\tau_k}^2) + A_{\tau_{k,w}} \psi(A_{\tau_{k,w}}^2) \leq \psi(n_{k+1}) \{A_{\tau_k} + A_{\tau_{k,w}}\} \leq 2\psi(n_{k+1}) A_{\tau_{k,w}}. \tag{61}$$

We evaluate

$$\mathcal{L}_{\tau_{k,w}}^{\gamma_k e^{-w}, k} := \int_{2/e}^1 \exp\left(u \gamma_k e^{-w} S_{\tau_{k,w}} - u^2 \gamma_k^2 e^{-2w} A_{\tau_{k,w}}^2 / 2\right) du$$

from above. Because  $u \gamma_k e^{-w} S_{\tau_{k,w}} - u^2 \gamma_k^2 e^{-2w} A_{\tau_{k,w}}^2 / 2$  is maximized at  $u = S_{\tau_{k,w}} / (\gamma_k e^{-w} A_{\tau_{k,w}}^2)$  and

$$\frac{S_{\tau_{k,w}}}{\gamma_k e^{-w} A_{\tau_{k,w}}^2} \leq \frac{2\psi(n_{k+1}) A_{\tau_{k,w}}}{(\psi(n_{k+1}) k^2 / \sqrt{n_{k+1}}) e^{-w} A_{\tau_{k,w}}^2} \leq \frac{2\sqrt{n_{k+1}}}{k^2 e^{-w} A_{\tau_{k,w}}} \leq \frac{2}{e^2} \leq \frac{2}{e},$$

the integrand in  $\mathcal{L}_{\tau_{k,w}}^{\gamma_k e^{-w}, k}$  is maximized at  $2/e$  and we have

$$\mathcal{L}_{\tau_{k,w}}^{\gamma_k e^{-w}, k} \leq \exp\left(\frac{2}{e} \gamma_k e^{-w} S_{\tau_{k,w}} - \frac{2\gamma_k^2 e^{-2w} A_{\tau_{k,w}}^2}{e^2}\right).$$

By (60) and (61), for sufficiently large  $k$ ,

$$\frac{2}{e} \gamma_k e^{-w} S_{\tau_{k,w}} - \frac{2\gamma_k^2 e^{-2w} A_{\tau_{k,w}}^2}{e^2} \leq \frac{4\gamma_k \psi(n_{k+1}) A_{\tau_{k,w}}}{e^{w+1}} - \frac{2\gamma_k^2 A_{\tau_{k,w}}^2}{e^{2(w+1)}}$$



$$\begin{aligned}
&= \frac{\psi(n_{k+1})^2 k^2 A_{\tau_{k,w}}}{\sqrt{n_{k+1}} e^w} \left( \frac{4}{e} - \frac{2k^2 A_{\tau_{k,w}}}{e^2 \sqrt{n_{k+1}} e^w} \right) \\
&\leq \frac{\psi(n_{k+1})^2 k^2 A_{\tau_{k,w}}}{\sqrt{n_{k+1}} e^w} \left( \frac{4}{e} - \frac{2}{e^2} \sqrt{e^4 - \frac{(1+\delta)k^4 n_k}{n_{k+1} e^{2w}}} \right) \\
&\leq -\psi(n_{k+1})^2 \frac{k^2}{\sqrt{n_{k+1}} e^w} \times \frac{\sqrt{n_{k+1}} e^{w+2}}{k^2} \times \frac{1}{2} \\
&= -\frac{e^2 \psi(n_{k+1})^2}{2}.
\end{aligned}$$

The last inequality holds because  $\lim_{k \rightarrow \infty} k^4 n_k / n_{k+1} = 0$  and  $4/e - 2 < -1/2$ . Hence  $\mathcal{L}_{\tau_{k,w}}^{\gamma_k e^{-w}, k} \rightarrow 0$  uniformly in  $1 \leq w \leq \lceil \ln k \rceil$ . This implies  $\limsup_{k \rightarrow \infty} \mathcal{M}_{\tau_{k+1} - \tau_k}^{\gamma_k, k} \leq 0$ .  $\square$

**Proposition 4.11.** *Let  $\omega \in \Omega_C$ . Suppose that  $v_k \leq \min(\tau_{k+1}, \sigma_{k,C})$  and*

$$-A_n \psi^U(A_n^2) \leq S_n, \quad \forall n \in [\tau_k, v_k].$$

*Then for sufficiently large  $k$*

$$\mathcal{N}_{v_k}^{\gamma_k, D} \geq \frac{\alpha}{2}.$$

*Proof.* As in the proof of the previous lemma, we denote  $t = n - \tau_k$ ,  $S_t = S_n - S_{\tau_k}$  and  $A_t^2 = A_n^2 - A_{\tau_k}^2$ . We distinguish two cases:

$$(a) A_{v_k}^2 \leq \frac{\psi(n_{k+1})^2}{c_1 \gamma^2}, \quad (b) \frac{\psi(n_{k+1})^2}{c_1 \gamma^2} < A_{v_k}^2 \leq A_{\tau_{k+1}}^2 - A_{\tau_k}^2.$$

For case (a), for sufficiently large  $k$  and for any  $\gamma \leq \gamma_k$ , as in (58),

$$\begin{aligned}
\gamma S_{v_k} &\leq \gamma (S_{v_k-1} + c_{v_k}) \leq \gamma \left( (1+\delta) A_{\tau_k} + \sqrt{A_{\tau_k}^2 + A_{v_k-1}^2} \right) \psi(n_{k+1}) + (1+\delta) C \frac{\sqrt{A_{\tau_k}^2 + A_{v_k-1}^2}}{\psi(A_{\tau_k}^2)^3} \\
&\leq \psi(n_{k+1})^2 \left( \frac{1}{\sqrt{c_1}} + \delta \right)
\end{aligned}$$

and

$$\psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} 2e^{(1+\delta)^5 e^6 C} e^{\gamma S_{v_k}} \rightarrow 0 \quad (k \rightarrow \infty).$$

Hence  $\mathcal{N}_{v_k}^{\gamma_k, D} \geq \alpha/2$  uniformly in  $\gamma \in [\gamma_k/k, \gamma_k]$ .

For case (b),  $S_{v_k}$  can be evaluated as

$$\begin{aligned}
S_{v_k} &\leq S_{v_k-1} + c_{v_k} \leq S_{v_k-1} + (1+\delta) C \frac{\sqrt{A_{\tau_k}^2 + A_{v_k-1}^2}}{\psi(A_{\tau_k}^2)^3} \\
&\leq \left( (1+\delta) A_{\tau_k} + \sqrt{A_{\tau_k}^2 + A_{v_k}^2} \right) \psi(n_{k+1}) + (1+\delta) C \frac{\sqrt{A_{\tau_k}^2 + A_{v_k}^2}}{\psi(A_{\tau_k}^2)^3}
\end{aligned}$$

$$\leq \left( (1 + \delta)A_{\tau_k} + \sqrt{A_{\tau_k}^2 + A_{\nu_k}^2} \left( 1 + \frac{(1 + \delta)C}{\psi(A_{\tau_k}^2)^3 \psi(n_{k+1})} \right) \right) \psi(n_{k+1})$$

by (51). Put

$$q_k^2 := \frac{A_{\tau_k}^2}{A_{\nu_k}^2} \leq \frac{c_1 \gamma_k^2}{\psi(n_{k+1})^2}, \quad s_k := \frac{(1 + \delta)C}{\psi(A_{\tau_k}^2)^3 \psi(n_{k+1})},$$

so that  $\lim_k q_k \psi(n_{k+1})^2 = 0$  and  $\lim_k s_k \psi(n_{k+1})^2 = 0$ . Then for sufficiently large  $k$

$$\begin{aligned} \frac{S_{\nu_k}^2}{2A_{\nu_k}^2} &\leq \left( (1 + \delta)^2 \frac{q_k^2}{2} + (1 + \delta)(1 + s_k)q_k \sqrt{1 + q_k^2} + (1 + s_k)^2 \left( \frac{1}{2} + \frac{q_k^2}{2} \right) \right) \psi(n_{k+1})^2 \\ &\leq \frac{\psi(n_{k+1})^2}{2} + \delta. \end{aligned}$$

Then

$$\psi(n_{k+1})e^{-\psi(n_{k+1})^2/2} \times 2e^{(1+\delta)^5 e^6 C} e^{S_{\nu_k}^2/(2A_{\nu_k}^2)} \frac{\sqrt{2\pi}}{\gamma A_{\nu_k}} \leq 2e^{(1+\delta)^5 e^6 C + \delta} \sqrt{2\pi c_1} e^\delta \leq \frac{D\alpha}{4}.$$

□

## 4.6 Dynamic strategy forcing the sharpness

Finally, we prove Proposition 4.1. We assume that by the validity result, Skeptic already employs a strategy forcing  $S_n \geq -A_n \psi^U(A_n^2)$  a.a. for  $\omega \in \Omega_C$ . In addition to this strategy, based on Proposition 4.10, consider the following strategy.

Start with initial capital  $\mathcal{K}_0 = \alpha$ .

Set  $k = 1$ .

Do the followings repeatedly:

1. Apply the strategy in Proposition 4.10 for  $n \in [\tau_k, \tau_{k+1}]$ .  
If  $\tau_{k+1} < \min(\sigma_{k,C}, \nu_k)$ , then go to 2. Otherwise go to 3.
2. Let  $k = k + 1$ . Go to 1.
3. Wait until  $\exists k'$  such that  $-\sqrt{\tau_{k'}} \psi^U(\tau_{k'}) \leq S_{\tau_{k'}} \leq \sqrt{\tau_{k'}} \psi(\tau_{k'})$ . Set  $k = k'$  and go to 1.

By this strategy Skeptic keeps his capital non-negative for every path  $\omega$ . For  $\omega \in \Omega_0$ ,  $\tau_k = \infty$  for some  $k$  and Skeptic stays in Step 1 forever. For  $\omega \in \Omega_{=\infty}$ , Step 3 is performed infinite number of times, but the overshoot of  $|x_n|$  in Step 3 does not make Skeptic bankrupt by Proposition 4.11. Now consider  $\omega \in \Omega_C$ . Since Skeptic already employs a strategy forcing  $S_n \geq -A_n \psi^U(A_n^2)$  a.a., the lower bound in (54) violated only finite number of times. By  $\omega \in \Omega_C$ ,  $n \geq \sigma_{k,C}$  is happens only finite number of times. Hence if  $S_n \leq A_n \psi(A_n^2)$  a.a., then Step 3 is performed only finite number of times and there exists  $k_0$  such that only Step 2 is repeated for all  $k \geq k_0$ . Now for each iteration of Step 2, Skeptic multiplies his capital at least by

$$1 + \frac{1 - \delta}{D} \lceil \ln k \rceil \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2}.$$

Then

$$\frac{1 - \delta}{D} \sum_{k=k_0}^{\infty} \lceil \ln k \rceil \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} \leq \prod_{k=k_0}^{\infty} \left( 1 + \frac{1 - \delta}{D} \lceil \ln k \rceil \psi(n_{k+1}) e^{-\psi(n_{k+1})^2/2} \right).$$

Since the left-hand side diverges to infinity, the above strategy forces the sharpness.

## References

- [1] Z. D. Bai. A theorem of Feller revisited. *Ann. Probab.*, 17(1):385–395, 1989.
- [2] I. Berkes, S. Hörmann, and M. Weber. Upper-lower class tests for weighted i.i.d. sequences and martingales. *J. Theoret. Probab.*, 23(2):428–446, 2010.
- [3] M. Csörgő, B. Szyszkowicz, and Q. Wang. Darling-Erdős theorem for self-normalized sums. *Ann. Probab.*, 31(2):676–692, 2003.
- [4] V. H. de la Peña, M. J. Klass, and T. L. Lai. Self-normalized processes: exponential inequalities, moment bounds and iterated logarithm laws. *Ann. Probab.*, 32(3A):1902–1933, 2004.
- [5] U. Einmahl and D. M. Mason. Some results on the almost sure behavior of martingales. In *Limit theorems in probability and statistics (Pécs, 1989)*, volume 57 of *Colloq. Math. Soc. János Bolyai*, pages 185–195. North-Holland, Amsterdam, 1990.
- [6] P. Erdős. On the law of the iterated logarithm. *Ann. of Math. (2)*, 43:419–436, 1942.
- [7] W. Feller. The general form of the so-called law of the iterated logarithm. *Trans. Amer. Math. Soc.*, 54:373–402, 1943.
- [8] W. Feller. The law of the iterated logarithm for identically distributed random variables. *Ann. of Math. (2)*, 47:631–638, 1946.
- [9] P. S. Griffin and J. D. Kuelbs. Some extensions of the LIL via self-normalizations. *Ann. Probab.*, 19(1):380–395, 1991.
- [10] N. C. Jain, K. Jogdeo, and W. F. Stout. Upper and lower functions for martingales and mixing processes. *Ann. Probab.*, 3:119–145, 1975.
- [11] A. Kolmogoroff. Über das Gesetz des iterierten Logarithmus. *Math. Ann.*, 101(1):126–135, 1929.
- [12] M. Kumon, A. Takemura, and K. Takeuchi. Capital process and optimality properties of a Bayesian skeptic in coin-tossing games. *Stoch. Anal. Appl.*, 26(6):1161–1180, 2008.
- [13] K. Miyabe and A. Takemura. Convergence of random series and the rate of convergence of the strong law of large numbers in game-theoretic probability. *Stochastic Process. Appl.*, 122(1):1–30, 2012.
- [14] K. Miyabe and A. Takemura. The law of the iterated logarithm in game-theoretic probability with quadratic and stronger hedges. *Stochastic Process. Appl.*, 123(8):3132–3152, 2013.
- [15] W. Philipp and W. F. Stout. Invariance principles for martingales and sums of independent random variables. *Math. Z.*, 192(2):253–264, 1986.
- [16] P. Révész. *Random Walk in Random and Non-Random Environments*. World Scientific Publishing Co. Pte. Ltd., Hackensack, NJ, third edition, 2013.
- [17] G. Shafer and V. Vovk. *Probability and Finance: It's only a game!* Wiley Series in Probability and Statistics. Financial Engineering Section. Wiley-Interscience, New York, 2001.

- [18] W. F. Stout. A martingale analogue of Kolmogorov's law of the iterated logarithm. *Z. Wahrscheinlichkeitstheorie und Verw. Gebiete*, 15:279–290, 1970.
- [19] V. Strassen. Almost sure behavior of sums of independent random variables and martingales. In *Proc. Fifth Berkeley Sympos. Math. Statist. and Probability (Berkeley, Calif., 1965/66), Vol. II: Contributions to Probability Theory, Part I*, pages 315–343. Univ. California Press, Berkeley, Calif., 1967.